

§ 3. Hamiltonian Dynamics:

Lecture 4

Definitions, Basic facts, Examples

10/05-2021

(M^{2n}, ω) symplectic

• $H: \mathbb{R} \times M \rightarrow \mathbb{R}$ — Hamiltonian
 t $H(t, \cdot) = H_t$

• Autonomous if ind of t : $H: M \rightarrow \mathbb{R}$

• often 1-periodic in t , $\mathbb{S}^1 = \mathbb{R}/\mathbb{Z}$
 $H: \mathbb{S}^1 \times M \rightarrow \mathbb{R}$ $H_{t+1} = H$

Def. • Hamiltonian v.f. generated by H :

$$i_{X_H} \omega = -dH \quad \exists! X_H \text{ — non-deg}$$

• $X_H \rightsquigarrow$ "time-dependent flow"
time dependent \nearrow
 $\varphi_H^t \leftarrow$ Hamiltonian flow generated by H
isotopy

Need not be defined for all t , and is not sometimes (collisions) but we will assume it's. (E.g. M is compact, etc)

Rem. Doing dynamics, usually interested in
in • φ_H^t , $t \in \mathbb{R}$, H autonomous
or φ_H^{kT} , $k \in \mathbb{N}$, H time-dependent

Examples

Ex 1. \mathbb{R}^{2n} , $\omega_{st} = dp \wedge dq = \sum dp_i \wedge dq_i$

$$\begin{cases} \dot{p} = -\frac{\partial H}{\partial q} \\ \dot{q} = \frac{\partial H}{\partial p} \end{cases} \Leftrightarrow X_H = -\frac{\partial H}{\partial q} \frac{\partial}{\partial p} + \frac{\partial H}{\partial p} \frac{\partial}{\partial q}$$

Checking $i_{X_H} \omega = -\frac{\partial H}{\partial q} dq - \frac{\partial H}{\partial p} dp = -dH$

Subexample $\mathbb{R}^{2n} = T^*\mathbb{R}^n = \mathbb{R}^n \times \mathbb{R}^n$

$H = \frac{1}{2m} \|p\|^2 + V(q) = \text{kinetic } p + \text{potential } q$

$$\begin{cases} \dot{p} = -\frac{\partial V}{\partial q} \Leftrightarrow \boxed{m\ddot{q} = -\nabla V} \leftarrow \text{cons force} \\ \text{Newton's eq} \end{cases}$$

$$\begin{cases} \dot{q} = \frac{1}{m} p \Leftrightarrow p = m\dot{q} \leftarrow \text{momentum} \end{cases}$$

2. Cotangent bundle $M = T^*Q$, ω_{st}

Fix R.M. on Q
 \langle, \rangle

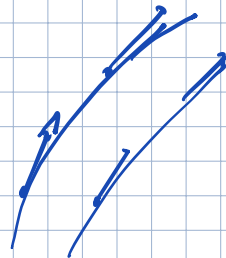
$TQ \leftrightarrow T^*Q$
 $\sigma \leftrightarrow \langle \sigma, \cdot \rangle$

$H = \frac{1}{2} \langle, \rangle : T^*Q \rightarrow \mathbb{R}$

$X_H = \text{geodesic spray}$

$\varphi_H^t = \text{geodesic flow}$

Describe.



Motion of a
free particle
on Q .

(17) on Q.

3. Twisted cotangent bundle

$$M = T^*Q, \quad \omega = \omega_{st} + \pi^*\sigma$$

$\downarrow \pi$
 $Q, \sigma, \quad d\sigma = 0$

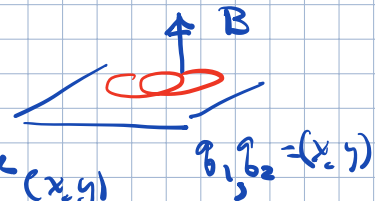
← magnetic field

$$H = \frac{1}{2} \langle \cdot, \cdot \rangle$$

The flow governs the motion of a charge on Q in magnetic field σ .

Subexample a) $Q = \mathbb{R}^2$, $\sigma = B dg_1 \wedge dg_2$

conf. space $\rightarrow (q_1, q_2)$



unit charge, unit mass
 $\vec{B} \perp (q_1, q_2)$ plane, charge in $\mathbb{R}^2 (x, y)$

H.E. $\Leftrightarrow \dot{\vec{q}} = B(q) J \dot{\vec{q}}; \quad J = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$

b) $Q = \mathbb{R}^3$, $\vec{B} = \text{v.f. on } \mathbb{R}^3 \leftarrow \text{magn. field}$

$$\sigma = \frac{1}{B} dg_1 \wedge dg_2 \wedge dg_3$$

$d\sigma = 0 \Leftrightarrow \text{div } B = 0$ } one of the Maxwell eq

H.E. $\Leftrightarrow \dot{\vec{q}} = \dot{\vec{q}} \times \vec{B}(q)$ ← Lorentz force

(unit charge, unit mass)

Energy and ω -conservation

Let φ_H^t be the flow of H_t .

Prop

essential

(a) Energy conservation

Assume that H is autonomous. Then

$$(\varphi_H^t)^* H = H : H(\varphi_H^t(p)) = \text{const } \forall p$$

(b) "Integral invariant": φ_H^t is symplectic

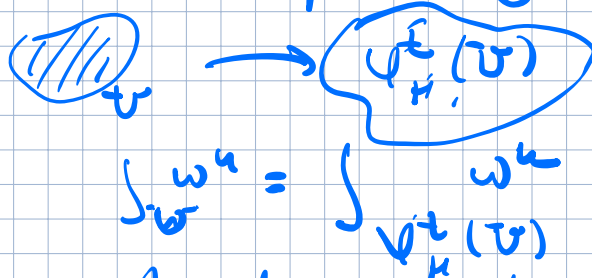
$$(\varphi_H^t)^* \omega = \omega$$

$$\varphi_H^t(\Sigma) = \Sigma'$$



$$\int_{\Sigma'} \omega = \int_{\Sigma} \omega$$

Cor. φ_H^t is vol. preserving: $(\varphi_H^t)^* \omega^k = \omega^k$



$$\int_U \omega^k = \int_{\varphi_H^t(U)} \omega^k$$

\Rightarrow these restrictions on dynamics

Pf. (a) $\frac{d}{dt} H(\varphi_H^t(p)) = (L_{X_H} H)(\varphi_H^t(p))$

$$= (i_{X_H} dH + d i_{X_H} H)(\dots)$$

(p. 11a)

$$= d(-i_{X_H} i_{X_H} \omega) = d\omega(X_H, X_H) = 0$$

HE: $i_{X_H} \omega = -dH$

$$\begin{aligned}
 (b) \quad \frac{d}{dt} (\psi_H^t)^* \omega &= (\psi_H^t)^* L_{X_H} \omega \quad (\text{see p 11a}) \\
 &= (\psi_H^t)^* \left(di_{X_H} \omega + i_{X_H} d\omega \right) \\
 &= (\psi_H^t)^* \underbrace{d(-dH)}_{HE} = 0
 \end{aligned}$$

Ex. Prove the Prop using a direct calculation in Darboux coordinates.

Con $\dim M = 2 \leftarrow$ surface
 H autonomous sol of ODE

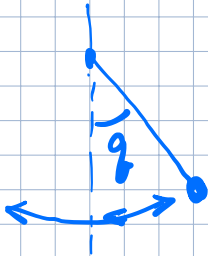
\Rightarrow integral curves of ψ_H^t (unparametrized)
 "ave" leaves $H = \text{const.}$
 "alg equations"

Prmk Newtonian mechanics:

- $\dot{q} = F(q)$ Energy cons $\Leftarrow \begin{cases} F \text{ is const.} \\ F = -\nabla V \leftarrow \text{indep of } t \end{cases}$
- $\dot{q}^0 = F(q, \dot{q})$ as in Lorentz
 Energy conservation $\Leftarrow F \perp \dot{q}$

Continuing Examples

4. Investigating the pendulum



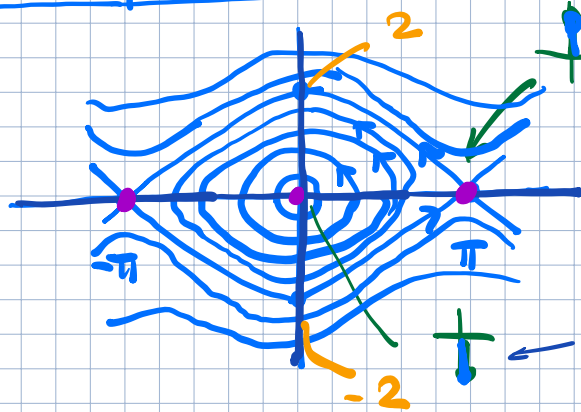
$$M = T^*S^1 = \mathbb{R} \times \begin{matrix} p \\ q \end{matrix} \leftarrow \text{mod } 2\pi$$

or $\mathbb{R}^2 = \mathbb{R} \times \mathbb{R}$

$$H = \frac{1}{2} |p|^2 - \cos q + 1$$

$$= \frac{1}{2} (p^2 + q^2) + \dots$$

Phase portrait:



← Harmonic oscillator
unstable
 $m=1, k=1$

$$\ddot{q} = -\sin q$$

Hooke's law:
 $\ddot{q} = -q$

Gives the behavior of integral curves up to parametrization

Further details - Ex (Not easy)

(a) $q \in (0, \pi)$ $H(0, q) = \cos q - 1 = h$

↑ integral curve through $(q, 0)$ is $\{H=h\}$
T(h) its period \int_x^x

• Show that T(h) monotone increasing function from 2π to ∞ as $h \rightarrow 2$

(21)

(Compare with the harmonic oscillator)
 $H = \frac{1}{2}(p^2 + q^2) \Leftarrow T = \text{const}$

- Find the Taylor exp of $T(h)$ at $h=0$
- (b) Consider $D\psi_H^T: T_x \mathbb{R}^2 \rightarrow \mathbb{R}^2$
Show that $D\psi_H^T = \begin{pmatrix} 1 & \lambda \\ 0 & 1 \end{pmatrix} \neq 0$
- (c) Find explicitly $\psi_H^t(0, 2)$
in elementary functions

Ex Hint to a): Area-period relation

$H: \mathbb{R}^2 \rightarrow \mathbb{R}$ proper

- $\{H \leq h\}$ connected
- h -regular

$$A(h) = \text{area of } \{H \leq h\} = \int_{H \leq h} \omega$$
$$T(h) = \text{period of } \{H = h\}$$

Show that $\frac{dA}{dh} = T(h)$

5. Positive Def quadratic Hom Lecture 5 10/07-2021

$$\mathbb{R}^{2n} = \mathbb{C}^n \quad z_j = p_j + iq_j, \quad z = (z_1, \dots, z_n)$$

$$H = \frac{1}{2} \sum \lambda_j (p_j^2 + q_j^2) = \frac{1}{2} \sum \lambda_j |z_j|^2$$

0 or just $\neq 0$

Describes n uncoupled oscillators
with frequencies λ_j

$E = \{H = h\}$ is an ellipsoid

- Find X_H and

show that $\varphi_H^t(z) = (e^{\lambda_1 t} z_1, \dots, e^{\lambda_n t} z_n)$

- "Coordinate axis" $(0, \dots, 0, z_j, 0, \dots, 0) \cap E$
are periodic orbits of φ_H^t with $T_j = \frac{2\pi}{\lambda_j}$

- Are there other periodic orbits?
(The answer depends on $(\lambda_1, \dots, \lambda_n)$.)

6. Linear HE

$$M = \mathbb{R}^{2n} \\ = \mathbb{C}^n$$

$$\omega = \omega_{st} = \text{dprdg}$$

matrix of ω :

$$J = \begin{pmatrix} 0 & -1 & & \\ 1 & 0 & & \\ & & \boxed{\begin{matrix} 0 & -1 \\ 1 & 0 \end{matrix}} & \\ 0 & & & \ddots \end{pmatrix}$$

← multipl. with
by i in $\mathbb{C}^n = \mathbb{R}^{2n}$

$H: \mathbb{R}^{2n} \rightarrow \mathbb{R}$ quadratic form

$$H(x) = \frac{1}{2} \langle Ax, x \rangle, \quad A^T = A$$

$$A = \nabla^2 H$$

HE: $\dot{x} = J \nabla H(x) = JA x = X_H(x)$

$$\Psi_H^t(x) = \exp(-tJA)x$$

Discusses a bit
Lie alg & Lie grps

$$\text{exp: } \mathfrak{sp}(2n) \rightarrow \text{Sp}(2n)$$

$$X_H \mapsto \Psi_H$$

$$\mathfrak{sp}(2n) = \text{lin Ham v.f.} \iff XJ + JX^T = 0$$

$$= \text{quadratic forms on } \mathbb{R}^{2n}$$

$$X \mapsto H = -\frac{1}{2} \langle JXx, x \rangle$$

(24)

Normal forms

Discuss linear alg

$Sp(2n)$ vs SL

• solving ODE's: $\dot{x} = Px$

Bring A to a Jordan form
to calculate $\exp(Pt)x$

$$PJP^T = J$$

$$P \in Sp(2n)$$

• symplectic normal forms are
more complicated

\mathbb{Q} symmetric

$$SQS^T \rightsquigarrow \text{diag}(1, \dots, 1, 0, \dots, 0)$$

$\cong GL(n)$

$$SQS^T \rightsquigarrow \text{diag}(\lambda_1, \dots, \lambda_n)$$

$\cong O(n)$

$SQS^T \rightsquigarrow$ much more complicated list

$\cong Sp(2n)$

Not just $\sum \lambda_j |z_j|^2$

Rmk. Euler $A > 0$, then it can be diagonalized